# Numerical investigation of heat transfer with natural convection in a regularly heated elliptical cylinder submerged in a square fence loaded with a nanofluid

Djedid Taloub<sup>1, 2</sup>, Adelkarim Bouras<sup>1</sup>, Zied Driss<sup>3</sup>

<sup>1</sup>Department of Physics, Faculty ofSciences, University Mohamed Boudiaf of M'sila, , Algeria

<sup>2</sup>Laboratory of Materials Physics and its Applications, University Mohamed BOUDIAF of M'sila,

Algeria

<sup>3</sup>Department of Mecanics, Electromechanical Systems Laboratory, (ENIS), University of Sfax, Tunisia

Received: August 13, 2021. Revised: November 9, 2021. Accepted: November 27, 2021. Published: December 16, 2021.

Abstract—During this first paper, numerical research from the natural convection of steady-state laminar heat transfer into a horizontal ring within a heated internal elliptical surface and a cold external square surface is presented. A Cu - water nanofluid, traverses this annular space. For different thermal Rayleigh numbers varying from 10<sup>3</sup> to 2.5x10<sup>5</sup> and different volume fractions from the nanoparticles. The arrangement from equations directing the problem was resolved numerically with the Fluent computational language founded on the finite volume approach. Based approaching the Boussinesq approach. The interior and exterior surfaces from the two cylinders are maintained at a fixed temperature. We investigated the impacts of various thermal Rayleigh numbers, the volume fraction from the nanoparticles, and the effect of the eccentricity of the internal cylinder on the natural convection. The results are shown within the figure of isocurrents, isotherms, and mean and local Nusselt numbers. The objective of this investigation is to examine the impact of different parameters on the heat transfer flow.

Keywords—a natural convection, Rayleigh numbers, nanofluid, volume fraction, square elliptic.

#### I. INTRODUCTION

THE necessity to enhance the thermal transfers of fluids has presented rise to the generation of a new class, called nanofluids. Holding nanoparticles suspended these solutions within a basis fluid. Gratitude to their thermic performanceenhancing properties, nanofluids can be applied in many fields such as the environment, production, cooling of electronic components, biology, medical diagnostics, water treatment, cooling in a heat engine Tzeng et al. [1], and energy storage. However, the main difficulty encountered with nanofluids is to disperse in a stable manner over time the nanoparticles in the base liquid because of their agglomeration, due either to gravity or to temperature-dependent precipitation. While the event of thermic convection in nanofluids, which are usually excellent conductors, both thermic and electrical, and into the attendance of a magnetic domain, two-volume forces control the nanofluid namely: the buoyance force and that of Lorentz. The end can cause magnetohydrodynamics (MHD).

The improvement from heat transfer by convection is the main object of several works, and to do so, a considerable number of investigations have carried out a throng of theoretical, experimental, and numerical essays relating to the representation of the aspects governing convection, the effect of the nature of the systems in which it takes place and the properties of the fluids involved, which are directly related to our study.

Mehrizi et al. [2] applied Boltzmann's approach to exploring the influence of nanoparticles on heat transfer by natural convection in a ring-shaped enclosure, which formed within a heated triangular internal surface and a circular external surface. The impact of the volume fraction from nanoparticles on improving heat transfer has been tested at various thermal Rayleigh numbers. In addition, the influence concerning diagonal, horizontal, and vertical eccentricities in different situations is tested at  $Ra = 10^4$ . These effects prove that the Nusselt number and the best flux rise with rising the fixed volume fraction. The mean Nusselt number rises as the internal

## INTERNATIONAL JOURNAL OF MATHEMATICAL MODELS AND METHODS IN APPLIED SCIENCES DOI: 10.46300/9101.2021.15.29

cylinder goes descending, but reduces as the position of the internal cylinder moves horizontally. Habibi et al. [3] presented a numerical investigation of natural convection and heat transfer from the copper (Cu) -water nanofluid interior a horizontal eccentric annulus. The internal and external cylinders are maintained at uniform temperatures. The influences of eccentricity, radius report, and volume fraction parameter of nanoparticles, thermal Rayleigh and Prandtl numbers on average Nusselt number, streamlines, and isotherms are examined. Izadi et al. [4] examined numerically the influence from the ratio of the Richardson numbers about the walls at the laminar combined convection of a nanofluid circulating within a uniformly heated annulus. The secondary currents and shape of dimensionless axial speed and dimensionless temperature are showed and explained. Matin et al. [5] examined numerically the combined convection and heat transfer about Al<sub>2</sub>O<sub>3</sub>-water nanofluid interior a horizontal eccentric annulus by turning on the internal cylinder. The interior and exterior cylinders are maintained at various fixed temperatures. The effects of eccentricity, Reynolds number, and volume fraction parameter from nanoparticles, thermal Rayleigh number on mean Nusselt number, and isotherms and isocurrent are reviewed. Dawood et al. [6] published a numerical simulation of increasing heat transfer employing nanofluids within an elliptical ring with a regular heat flow boundary state. Four different types of nanofluids of various sizes from nanoparticles were used. They found that glycerin exhibited the best improvement in heat transfer related before the other base fluids were tested. Moghari et al. [7] have studied numerically the pressure loss and thermal properties of the flow of nanofluid Al<sub>2</sub>O<sub>3</sub> water in horizontal rings. Each single-phase fluid method was selected for the modeling of nanofluids. The effects of some important parameters studied and discussed in detail. They have observed that there are significant increases in pressure loss and in elevating energy, which is not acceptable. Seyyedi et al. [8] simulated heat transfer by natural convection under uniform heat flow into an annular enclosure loaded by nanofluids. They used the finite element approach based on the control volume. To evaluate the impact of thermic conductivity and viscidity from the nanofluid they used the Maxwell-Garnetts (MG) and Brinkman models. These results explicate that increasing the aspect report raises the values of the mean Nusselt number. Togun et al. [9] performed a collection from subjects about forced, natural, and combined heat transfer on fluid and nanofluid flow within an annular passing by various fluid and boundary states. The influence concerning eccentricity into the plane and perpendicular directions on the rate of heat transfer in the largest digital and empirical studies for plane and perpendicular annular passages investigated. The influences of Darcy, Prandtl number, Reynolds number, Grashof, and thermal Rayleigh numbers on heat transfer are investigated. The impact of the nanofluid on the further improvement of heat transfer within an annular canal is done. Garoosi et al. [10] performed numerical research at the heat transfer by

natural convection from nanofluid inside a rectangular cavity including different pairs from heaters and icehouse. The surfaces from the cavity are protected and many pairs of heating and cooling devices with isothermal surfaces are placed inside the cavity. The influences of many parameters about the heat transfer flow and the dispersion of nanoparticles are investigated. These results designate that the location of the coolers has several important impacts on the heat transfer rate. They noticed that there is an optimum volume fraction of nanoparticles by any thermal Rayleigh number into which the best heat transfer flow can be achieved. Arbaban et al. [11] improved the heat transfer into horizontal concentric rings with eight radial flippers connected to the internal cylinder using nanofluids.

These results confirm that the mean Nusselt number rises with the rising volume fraction and thermic conductivity from nanoparticles. In addition, it is mentioned that the mean Nusselt number from Cu-water nanofluid is the most important with the nanofluids. Boulahia et al. [12] numerically examined the performance from natural and combined convection from Cu-water nanofluid in square forms including circular warming and cooling cylinders inside. The effects of the position and size from circular shapes at the flow of heat transfer are studied. The results confirm that the optimal heat transfer is achieved by locating the circular form near the floor wall, and the rate concerning heat transfer rises with the change in guidance from the pair of circular shapes from the horizontal directorate to the vertical directorate. Husain et al. [13] presented a numerical and experimental examination concerning the transitory natural convection flux about distilled water into a narrow perpendicular ring with a high aspect report. Using ANSYS Fluent a CFD model is developed and solved. The analysis was carried out on a heat facts scale of 300 W to 1800 W. They observed that by an improvement in heat facts, a decrease in the transitory time. Hamid et al. [14] studied the increased warmth transfer and hydromagnetic flux of water-founded carbon nanotubes in a partly heated square fin-formed cavity. Inside the cavity, a fine heated bar is fixed as a heat source. In the calculation, the numerical finite element approach is performed utilizing the Galerkin method. These results illustrated that local Nusselt numbers are enhanced with injecting both a water-based solid carbon volume fraction and radiation influences, while the Nusselt number is maximal in the angles. Lakshmi et al. [15] analytically investigated the natural convection from watercopper nanoliquids into saturated cylindric porous rings. The internal and external perpendicular radial surfaces are subjected to regular mass and heat fluxes and outflows, respectively. The thermophysical characteristics of the saturated nanoliquid porous means are patterns utilizing phenomenological rules and mixing hypotheses. The influence of many parameters about velocity, temperature, and heat transport is found. The highest heat transport is obtained within a simple cylindrical ring contrasted to the rectangle and high circular rings. Increasing the radius from the internal

cylinder is to lowers the heat transport. Mebarek et al. [16] numerically studied the convective heat transfer of Titania nanofluids with various basis fluids into a cylindrical ring by a discrete warmth source. A computer code is a developed basis on the finite volume approach coupled with the SIMPLER algorithm. The influences from these parameters on the local Nusselt number are investigated. Siddheshwar et al. [17] analytically studied the linear and nonlinear Darcy-Bénard convection about Newtonian nonliquids and Newtonian liquids into cylindrical and cylindrical annuli. They studied the effect of heat transport and the concentric addition from a cylinder in the circular cylinder. The two-phase representation is utilized within the event of nanoliquids. They noticed that the heat transportation is highest in the event from a cylindrical ring accompanied on such cylindrical and square enclosures. Mebarek et al. [18] numerically analyzed the effects from the position of a heat fount at the floating convection of nanofluids within an annular system. They examined five various heat fount locations alongside the internal cylinder from the ring. Their goal is to recognize the optimal location from the warmth fount to maximize or minimize heat transportation to many Rayleigh number valors and various volume fractions from the nanoparticle. These results show that the position from the heat fount has a deep influence on the flow motif. The volume fraction from nanoparticles more commands heat transportation into the ring geometry. Keerthi et al. [19, 20] analyzed the impact of a non-uniform warm state at floating convective transport from nanofluids inside a cylindrical ring formed by two vertical coaxial cylinders. Numerical simulation results show the bicellular flow model to the two non-uniform thermic states across any Rayleigh number ranges. Lakshmi et al. [21-24] examined analytically the natural convection of nanoliquids into saturated porous cylindrical rings with the inner cylinder under regular mass and heat flow, and the outer cylinder under heat flow. The Darcy pattern and the changed account of the two-phase Buongiorno pattern are utilized. Of the three examples of rings, estimated, shallow rings present the most suitable heat transportation and high rings present the most critical production. The appearance concerning a dilute concentration from nanoparticles significantly improves the warmth transportation within the system. They concluded that a simple cylindrical ring soaked by water and AA7075 lamellar mixture nanoparticles is properly satisfied to heat transfer due to its large efficient thermic conductivity compared to other formed nanoparticles and that a large quadrangular enclosure soaked with water strongly sequences to heat stocking applications. Sankar et al. [25] examined the effects from heat transport conjugated into nanofluids by many nanoparticles so as alumina, titanium oxide, or copper within a closed annular geometry and formed by an internal cylinder having a regular large temperature an outer cylindrical tubing by a fixed below temperature and thermic insulated top and bottom surfaces for a broad Rayleigh number spectrum. These results show that copper-based nanofuids produce more important heat transportation between different nanoparticles. Sultan et al. [26] studied numerical the heat transfer by natural convection using various nanofluids (Ag, Cu, and TiO<sub>2</sub>) in a horizontal three-dimensional flow and an inclined ring with two heated flippers are fixed on the internal cylinder while the outer cylinder is adiabatic. These numerical effects confirm that while the volume fraction rises, the warmth transfer improved to any valors about the Rayleigh number. The improvement is most important in large Rayleigh numbers. Roy [27] examined the buoyant motion of nanofluids in annular spaces within a rectangular geometry and three inside geometries: elliptical, circular, or rectangular cylinder. He noted that the interior forms have a deep impact on heat transfer flows compared to a square geometry.

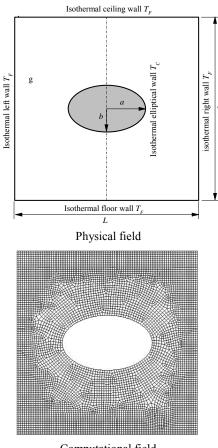
We conducted a literature search on the associated floatable convection of basic fluids and nanofluids under many geometries. Following detailed analytical and empirical studies, we remarked that the buoyancy movement and heat transfer behavior from nanofluids interior the closed square geometry have not been performed numerically in detail until now to analyze the impacts of several key parameters about the thermal behavior from nanofluids. The work we develop in this manuscript concerns the numerical research from warmth transfer by nanofluids: two coaxial cylinders, a cold external square cylinder, and a heated interior elliptical cylinder. Numerical simulations have been performed to predict the effect of Rayleigh numbers, the solid volume fraction of nanoparticles, and inner cylinder eccentricity at the flow field and structure, and heat transfer.

#### II. DETERMINING EQUATIONS TO LAMINAR NANOFLUIDS

#### A. Physical Models

Figure 1 shows the physical representation from the current work. The treated problem is an elliptical cylinder with an eccentricity  $e_1$  (= 0.8), located inside a square cylinder enclosure characterized by (*L*). The surfaces from the outer square fence have been maintained in a fixed low-temperature  $T_F$ , and the inner ellipse surface is kept in a fixed high-temperature  $T_C$ . During this work, the thermal Rayleigh's number,  $Ra_t$ , ranges from  $10^3$  to  $2.5 \times 10^5$ . The problem study is 2D, the flow is convection natural laminar of nanofluid (Cuwater). The nanofluid is considered incompressible and Newtonian by negligible viscous diffusion and pressure working. The solid and fluid phases are into thermic stability by the same velocity of flow and the Boussinesq approach used.

Volume 15, 2021



Computational field

Fig. 1 Physical and computational field

#### B. Mathematical model

The treatment of the considered physical problem requires the application of governing equations deduced from the classical principles of conservation of mass, momentum, and energy.

The determining equations converted into dimensionless models below the non-dimensional variables:

$$X = \frac{x}{L}, Y = \frac{y}{L}, U = \frac{u}{\binom{\alpha_f}{L}}, V = \frac{v}{\binom{\alpha_f}{L}}, \theta = \frac{T - T_c}{T_h - T_c}, P = \frac{p}{\rho(\frac{\alpha_f}{L})^2}, Pr = \frac{\theta_f}{\alpha_f}, Ra_t = \frac{g\beta_f (T_h - T_c)D_h^3}{\alpha_f \theta_f}$$

By introducing the dimensionless quantities into the conservation equations of mass, motion and energy, we obtain respectively:

$$\frac{\partial U}{\partial x} + \frac{\partial V}{\partial y} = 0 \tag{1}$$

$$U\frac{\partial U}{\partial x} + V\frac{\partial U}{\partial y} = -\frac{\partial P}{\partial x} + \frac{\mu_{\rm nf}}{\rho_{\rm nf}\,\alpha_f} \left(\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2}\right) \tag{2}$$

$$U\frac{\partial V}{\partial x} + V\frac{\partial V}{\partial y} = -\frac{\partial P}{\partial y} + \frac{\mu_{\rm nf}}{\rho_{\rm nf} \alpha_f} \left(\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2}\right) + \frac{(\rho \beta)_{\rm nf}}{\rho_{\rm nf} \beta_f} Ra_t Pr\theta$$
(3)

$$U\frac{\partial\theta}{\partial x} + V\frac{\partial\theta}{\partial Y} = \frac{\alpha_{nf}}{\alpha_f} \left( \frac{\partial^2\theta}{\partial x^2} + \frac{\partial^2\theta}{\partial Y^2} \right)$$
(4)

#### C. Estimation from the Nusselt number

The local Nusselt number adjacent to the outer and interior cylinders and the mean Nusselt number can be estimated as follows [28, 29]:

$$Nu_{h} = -\frac{k_{nf}}{k_{f}} \left(\frac{\partial T}{\partial Y}\right)_{Y=0}$$
(5)

$$Nu_{c} = -\frac{\kappa_{\text{nff}}}{k_{f}} \left( \frac{\partial I}{\partial Y} \right)_{Y=1}$$
(6)

$$\overline{Nu}_{avg} = \frac{1}{p_h} \int_0^{p_h} Nu_x \, dx \tag{7}$$

Table 1 Thermophysical characteristics from the base fluid and Cu nanoparticles [30]

	$\rho(\frac{Kg}{m^3})$	$C_p\left(\frac{J}{kg.K}\right)$	$K\left(\frac{W}{m.K}\right)$	$\beta\left(\frac{1}{K}\right)$	μ (Ps.s)
Water	997.1	4179	0.613	21x10 <sup>-5</sup>	0.000891
Cu	8933	385	401	1.67x10 <sup>-5</sup>	-

#### D. Nanofluid characteristics

The thermophysical characteristics of base fluid (water) and nanoparticles (copper) are presented in table 1. The characteristics of the nanofluids can be get from the properties of the basic fluid and the nanoparticles. The heat capacity and density of the nanofluids are estimated based on the references of Ramiar et al. [32] and Khanafer and Vafai [33].

$$(\rho)_{\rm nf} = \phi \rho_p + (1 - \phi)\rho_f \tag{8}$$

$$\left(\rho C_p\right)_{\rm nf} = \phi \left(\rho C_p\right)_p + (1-\phi) \left(\rho C_p\right)_f \tag{9}$$

The thermic diffusivity from the nanofluid is

$$(\alpha)_{\rm nf} = \frac{\kappa_{\rm nf}}{(\rho c_p)_{\rm nf}} \tag{10}$$

The thermal expansion coefficient of the nanofluid calculated by

$$(\rho \beta)_{nf} = \phi (\rho \beta)_p + (1 - \phi)(\rho \beta)_f \tag{11}$$

The efficient dynamic viscosity from the nanofluid determined by Brinkman [34] as

$$(\mu)_{nf} = \frac{\mu_f}{(1-\phi)^{2.5}} \tag{12}$$

Concerning effective thermic conductivity from the nanofluid, Maxwell [35] presented a model to base density mixing by micron-sized spherical particles, which is

$$K_{nf} = K_f \frac{(\kappa_p + 2\kappa_f) - 2\phi(\kappa_f - \kappa_p)}{(\kappa_p + 2\kappa_f) + \phi(\kappa_f - \kappa_p)}$$
(13)

#### E. Boundary condition

Solving the system of equations got previously requires the incorporating of boundary conditions to any dependant variable. The temperature conditions are identified on the cylinders. The temperature from the external cylinder is uniform and equal to  $T_f$ . The hot part of the inner cylinder has a constant temperature equal to  $T_c$ .

These different boundary conditions in dimensional form can be recapitulated as:

The initial conditions are:

$$U = V = 0 \tag{14}$$

$$\theta(X,Y) = 0 \tag{15}$$

Moreover, the boundary limitations about the problem are:

Inner cylinder (ellipse)

$$U = V = 0 \tag{16}$$

$$\theta(X,Y) = 1 \tag{17}$$

External cylinder (square)

$$U=V=0 \tag{18}$$

$$\theta(X,Y) = 0 \tag{19}$$

#### F. Meshing choice

The impact of the size and number from nodes at the resolution represented on the warmth transfer to the "heated" active party from the cylinder explained with the average Nusselt number in Figure 2. An irregular mesh close to the walls was utilized to resolve more exactly the physical phenoms present in special within the boundary layer zone designated with the presence of large gradients into the parietal regions.

To get a mesh-unattached solution, a grill refinement investigation is presented to a square fence by an isothermal internal ellipse at  $Ra_t = 10^6$  (figure 2). In this paper, eight combinations ( $80 \times 80$ ,  $90 \times 90$ ,  $100 \times 90$ ,  $110 \times 90$ ,  $115 \times 90$ ,  $120 \times 90$ , and  $100 \times 100$ ) of check volumes are utilized to examine the influence of grill size at the precision of foretold results. Figure 2 presents the convergence from the mean Nusselt number, in the heated inner cylinder level by grill refinement. It is remarked that mesh independence is reached by combining control volumes ( $100 \times 100$ ) in which the mean Nusselt number does not vary importantly with the finer mesh refinement. The accord was determined to be excellent, which corresponds to the current calculation indirectly.

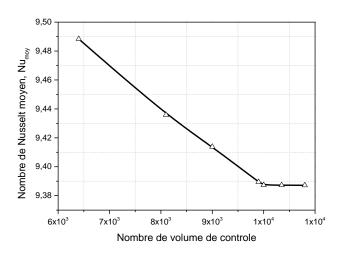


Fig. 2 Convergence from the mean Nusselt number alongside the hot interior cylinder with grill refinement to  $Ra_i = 10^6$ 

#### III. VALIDATION OF RESULTS

To verify the precision of the numerical results performed in this paper, validation of the simulation effected employing into statement any numerical investigations available in the research.

The phenom of heat transfer with natural convection into an external closed square fence at low temperature and an internal elliptical cylinder at high temperature has been treated. The average Nusselt numbers evaluated for the examination are related to the reference conditions of Kim et al. [31] as presented within table 2. Added validation was carried out employing the existing mathematical algorithm to study the very problem as that analyzed by [31] utilizing the identical flow provisions and geometries, which were related to laminar natural convection heat transfer utilizing the same boundary provisions, but the numerical scheme was different. The comparison is made utilizing the next dimensionless parameters: Pr = 0.7,  $Ra_t = 10^5 \cdot 10^6$ . The best arrangement was obtained within Kim et al. [31] and the current numerical simulation scheme to the temperature contours and streamlines inside the squared enclosure by the internal cylinder, as exposed in figure 3. Those validations give confidence within the current mathematical model employed.

Table 2 Comparison of the current mean Nusselt number by preceding numerical investigations.

Rat	Mean Nusselt number about the heated cylinder		Erreur
	Now investigate	Kim et al. [31]	- (%)
104	3.461	3.414	1,3767
10 <sup>5</sup>	5.0592	5.1385	1,5432

-

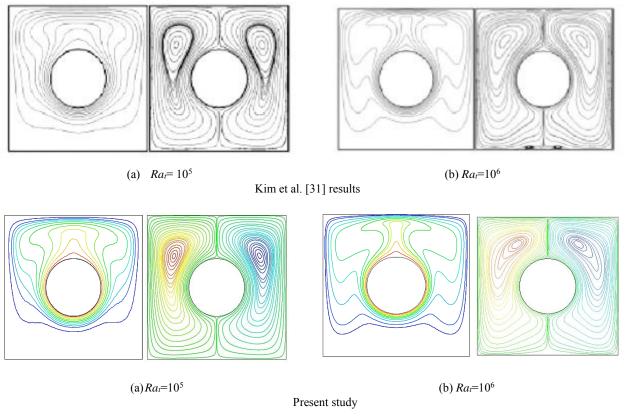


Fig. 3 Comparison of streamlines and temperature contours within the existing performance and such of Kim et al. [31]; concerning two various thermal Rayleigh numbers of (a) 10<sup>5</sup> and (b) 10<sup>6</sup>

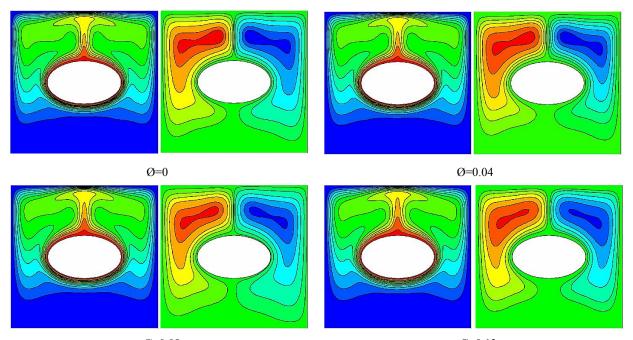
#### IV. RESULTS AND ANALYSIS

Within this document, we have offered numerical simulations about natural convection in a closed annular system, this system comprising two confocal cylinders an outer square and an inner ellipse, which controlled at uniform imposed temperatures. The system of equations ruling the problem is resolved numerically utilizing a finite volume approach based on the SIMPLE algorithm. The developed model is first validated from the numerical results. The thermal and dynamic current field is shown. The estimates are carried for thermal Rayleigh numbers of 10<sup>3</sup>, 10<sup>4</sup>, 10<sup>5</sup>, and 2.5x10<sup>5</sup>, and volume fraction of nanoparticles  $0 \le \emptyset \le 0.12$ . We examine the annular space defined by the eccentricity of the inner elliptical cylinder ( $e_1 = 0.8$ ), and the outer cylinder is a square cylinder designated by (L). This annular space is closed which contains a Cu-water nanofluid. The thermophysical characteristics of the copper nanoparticles and the base fluid are presented in table 1. The external and the internal wall create a temperature gradient.

Figure 4 plotted isocurrent and isotherms for various volume fraction estimations of nanoparticles when the thermic Rayleigh number equals 10<sup>5</sup>.

It is noted that the isotherms and the streamlines are symmetrical with respect to the fictitious vertical median plane. This figure show that the flow regime is single-cell. On the left side of the plane of symmetry, the flow rotates counterclockwise and at the right lateral, it is within the opposite sens (the particles of the fluid movement upwards below the action from gravity forces along the internal hot wall also then descend to the vicinity of the cold wall of the external cylinder).

At  $Ra_t = 10^5$ , the isotherms grow and ultimately take on the form of a mushroom. The temperature distribution decreases from hot cylinder to cold cylinder. The direction of the deformation of the isotherms corresponds to the path of rotation of the isocurrent. In laminar mode, it can be said that, under the action of the movement of the particles which take off from the hot wall at the level of the axis of symmetry, the isotherms "arch" and move away from the wall at this point. The preferences from the streamlines raise, which confirms that the convection is increasing. In fact, the existence of nanoparticles causes an intensification of isotherms near the hot surface, which indicates an enhancement in the rate of heat transfer.



 $\emptyset$ =0.08  $\emptyset$ =0.12 Fig. 4 Isocurrent (right) and Isotherms (left) concerning various values from volume fraction when,  $Ra = 10^5$ , and  $e_1 = 0.8$ 

Figure 5 presents the contours about the temperature and streamlines when the volume fraction values of the nanofluids are zero ( $\emptyset = 0$ ) for  $Ra_t = 10^3$ ,  $10^4$ ,  $10^5$ , and  $2.5 \times 10^5$ . The figure notes that the heated water next to the hot isothermal elliptical surface moves upward and then hits the square cylinder isothermal cold surface. Next striking the cold cylinder, it reverses direction creating a symmetrical rotating vortex.

Concerning isotherms, at thermal Rayleigh numbers are between  $10^3$  and  $10^4$ , the isotherms are generally identical, and the heat is transported by conduction. At  $Ra_t$  (=  $10^5$  and  $2.5 \times 10^5$ ), the temperature shapes and isocurrent enhance extra confusing, and the vortex intensity raises while the Rayleigh number rises. In that case, the flow highly strikes the ceiling from the square cylinder, which conducts the production of a thin thermic layer. The form of the rotating spirals increases, the isotherms go up, and heat is transferred by convection.

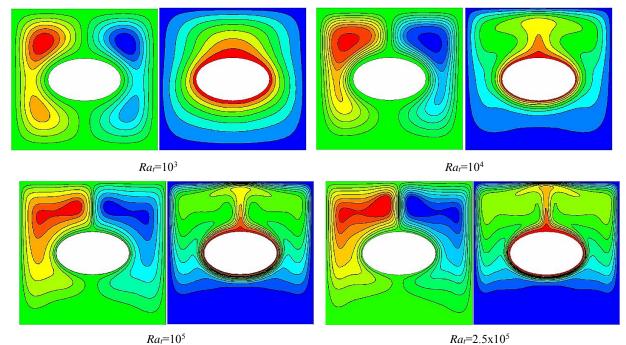


Fig. 5 Isotherms (right) and isocurrents (left) for several values about the thermal Rayleigh number,  $Ra_t$  at  $\emptyset=0$ , and  $e_1 = 0.8$ 

The evolution from the average Nusselt numbers by the nanoparticle volume fraction for eccentricity about the internal ellipse ( $e_1$ ) is presented in figure (6). It is noted that as an adjusted value concerning the thermal Rayleigh number, the mean Nusselt number rises linearly with increasing Ø. The slope of the line is increased when the thermal Rayleigh number and the volume fraction from the nanoparticles increase at this eccentricity, which leads to an increase in the heat flux strongly depending on the thermal Rayleigh number.

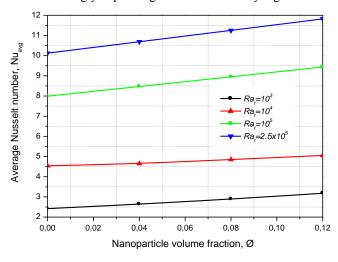


Fig. 6 Evolution of the mean Nusselt number by the nanoparticles volume fraction for various thermal Rayleigh number and  $e_1 = 0.8$ 

Figure. 7 presents the evolution from the mean Nusselt number in the function from the thermal Rayleigh number concerning different values about volume fraction from nanoparticles where the eccentricity  $e_1 = 0.8$  of the internal surface. It is apparent that with the rise in the volume fraction about nanoparticles, the efficient thermic conductivity from the nanofluid raises, which appears to enhance thermic transportation from the fluid inside the channel and better heat convection.

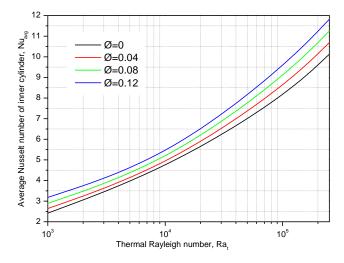


Fig. 7 Evolution from the average Nusselt number round an internal surface by the thermal Rayleigh to various values about the nanoparticle volume fraction and  $e_1 = 0.8$ 

The changes from the local Nusselt number by way alongside the area about the square forms (A, B, C, and D) and by several values about the nanoparticle volume fraction at thermal Rayleigh 10<sup>5</sup> are explained within the figure (8). The local Nusselt numeral evolutions are superimposed upon each other, and the local Nusselt number values rise by the rise in the value of the volume fraction from the nanoparticles. At thermal Rayleigh number 10<sup>5</sup>, as shown under figure (8), the local Nusselt number values increase due to the important influence of thermic convection and nanoparticle volume fraction values. It starts from the highest value in the summit region of the internal surface due to the temperature source effect of the inner cylinder and the presence of small vortices at the internal cylinder and goes down alongside the area of the fence before it approaches the least value near the point B. Next, the local Nusselt number gradually rises and next reductions until it arrival position C. Then, the local Nusselt number leaps directly from the area of point C to point D, where the values are almost zero it is the lower part of the enclosure where the fluid is practically stationary.

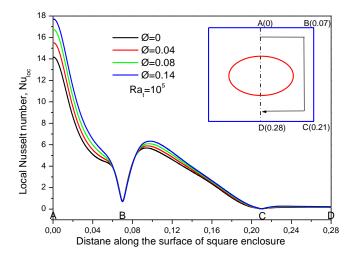


Fig. 8 Evolution of the local Nusselt number alongside the area (A– B–C–D) from the squarish fence at various values about the nanoparticle volume fraction for  $Ra_t = 10^5$  and  $e_1 = 0.8$ .

By fixing the dimensions of the external cylinder and varying the dimensions of the internal elliptical cylinder el and a given thermal Rayleigh. Figure (9) shows the development from the local Nusselt number alongside the area (A-B-C-D) of the squarish fence for diverse values of the eccentricity when  $Ra_t = 10^5$  and  $\emptyset = 0$ . We noticed that for the valor of thermal Rayleigh number considered shows that when  $e_1 = 0.6$  the liberate interspace between the two swap areas is short, it thus supports the conductive change, but by the increase about  $e_1$ , this liberate space will widen to support the change by convection which rests relatively weak for numbers of thermal Rayleigh weak, but this transfer by convection becomes more important for  $Ra_t = 10^5$ . The increase in the value of the inner cylinder eccentricity  $e_1$  thus leads to an increase in the liberate interspace between the two interchange surfaces, which results in an intension of the natural convection, for example for a thermal Rayleigh number ( $Ra_{l} = 10^{5}$ ), the figure shows no bifurcation for  $e_{l} = 0.7$  and  $e_{l} = 0.8$ , the evolutions of local Nusselt number are nearly identical and concentric trajectories, for example for  $e_{l} = 0.6$ , the figure illustrates the existence of a bifurcation and the presence of a bifurcation presenting rise to further cells rotating in the inverse direction of the adjacent cells and three feathers in the isotherms.

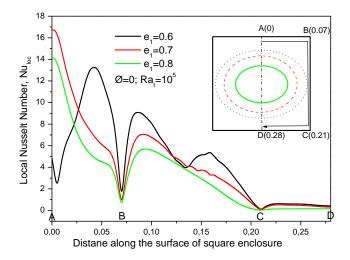


Fig. 9 Local Nusselt number evolution along the distance (A-B-C-D) of the squarish fence for various values of the eccentricity when  $Ra_t = 10^5$  and  $\emptyset = 0$ .

#### V. CONCLUSION

The enclosures delimited by two cylinders are particularly interesting because they allow a number of configurations and lead to original results, which show in particular that: If the enclosure opposes to the fluid movement sufficiently pronounced constrictions, it is possible to observe multicellular flows even at low values from the thermal Rayleigh number. In this case, the rise concerning the thermal Rayleigh number, i.e. the enhancement of the convection natural, leads to a raise of the vortices.

A numerical investigation from natural convection within a squared fence isothermal warmed by an internal elliptical surface with a copper–water-based nanofluid has been performed. The two-dimensional conservation equations of energy, mass, and momentum including the Boussinesq approach were determined utilizing the finite volume technique. The determining parameters are  $10^{3} \le Ra_{t} \le 2.5 \times 10^{5}$ ,  $0 \le \emptyset \le 0.12$ , and  $0.6 \le e_{I} \le 0.8$  within viewpoint about the obtained results, the conclusions obtained were summarized as follows:

• To a low thermal Rayleigh number, the conduction process essentially controls the heat transport within the annulus. During the thermal Rayleigh number rises, the role of convection becomes dominant. The isotherms grow more confounded and thermal convection is commanding.

• The contours of the streamlines and the isotherms are symmetrical to the perpendicular line for all situations.

• In high thermal Rayleigh numbers, a thin thermal

boundary layer is illustrated at the flow that heavily strikes the ceiling and lower from the outer cylinder.

• The addition of copper nanoparticles produced an important enhancement within the flow from heat transfer, so the heat transfer rises with the rise in the volume fraction of the nanoparticles.

• The local and mean Nusselt number from a nanofluid are enhanced by enhancing the volume fraction of the nanoparticles.

• The space between the two cylinders has a major significant impact on improving the rate of heat transfer when the convective mode dominates.

#### References

- S.C. Tzeng, C.W. Lin, K. Huang, "Heat transfer enhancement of nanofluids in rotary blade coupling of fourwheel-drive vehicles," Acta Mechanica., vol. 179, pp. 11-23, 2005.
- [2] A. A.Mehrizi, M.Farhadi, and S. Shayamehr, "Natural Convection Flow of Cu–Water Nanofluid in Horizontal Cylindrical Annuli with Inner Triangular Cylinder Using Lattice Boltzmann Method," Int. Commun. Heat Mass Transfer., vol. 44, pp. 147–156, 2013.
- [3] M. Habibi Matin, and I. Pop, "Natural Convection Flow and Heat Transfer in an Eccentric Annulus Filled by Copper Nanofluid," Int. J. Heat Mass Transfer, vol. 61, pp. 353–364, 2013.
- [4] M. Izadi, M. M. Shahmardan, A. Behzadmehr, "Richardson Number Ratio Effect on Laminar Mixed Convection of a Nanofluid Flow in an Annulus," International Journal for Computational Methods in Engineering Science and Mechanics, vol. 14, no. 4, pp. 304-316, 2013.
- [5] Meisam Habibi Matin and Ioan Pop, "Numerical Study of Mixed Convection Heat Transfer of a Nanofluid in an Eccentric Annulus," Numerical Heat Transfer, Part A: Applications, An I. J. of Computation and Methodology, vol. 65, no. 1, pp. 84-105, 2014.
- [6] H. Dawood, H. Mohammed, and K. Munisamy, "Heat Transfer Augmentation Using Nanofluids in an Elliptic Annulus With Constant Heat Flux Boundary Condition," Case Stud. Therm. Eng, vol. 4, pp. 32–41, 2014.
- [7] R. M. Moghari, F. Talebi, R. Rafee, and M. Shariat, "Numerical Study of Pressure Drop and Thermal Characteristics of Al<sub>2</sub>O<sub>3</sub>–Water Nanofluid Flow in Horizontal Annuli," Heat Transfer Eng, vol. 36, no. 2, pp. 166–177, 2014.
- [8] S. Seyyedi, M. Dayyan, S. Soleimani, and E. Ghasemi, "Natural Convection Heat Transfer Under Constant Heat Flux Wall in a Nanofluid Filled Annulus Enclosure," Ain Shams Eng. J, vol. 6, no. 1, pp. 267–280, 2015.
- [9] Hussein Togun, Tuqa Abdurazzaq, S.N. Kazi, Emad Sadeghinezhad, "A review of studies on forced, natural and mixed heat transfer to fluid and nanofluid flow in an annular passage," Renewable and Sustainable Energy Reviews, vol. 39, pp. 835-856, 2014.

- [10] F. Garoosi, L. Jahanshaloo, M. M. Rashidi, A. Badakhsh, and M. E. Ali, "Numerical Simulation of Natural Convection of the Nanofluid in Heat Exchangers Using a Buongiorno Model," Appl. Math. Comput vol. 254, pp. 183–203, 2015.
- [11] M. Arbaban, and M. Salimpour, "Enhancement of Laminar Natural Convective Heat Transfer in Concentric Annuli With Radial Fins Using Nanofluids," Heat Mass Transfer, vol. 51, no. 3, pp. 353–362, 2015.
- [12] Zoubair Boulahia, Abderrahim Wakif, J. Ali Chamkha, Rachid Sehaqui, "Numerical study of natural and mixed convection in a square cavity filled by a Cuwater nanofluid with circular heating and cooling cylinders," Mechanics and Industry, vol. 18, no. 5, pp. 1, 2017.
- [13] S. Husain, & M. A. Siddiqui, "Experimental and numerical analysis of transient natural convection of water in a high aspect ratio narrow vertical annulus," Prog. Nucl. Energy, vol. 106, pp. 1–10, 2018.
- [14] M. Hamid, Z. H. Khan, W. A. Khan, R. U. Haq, "Natural convection of water-based carbon nanotubes in a partially heated rectangular fin-shaped cavity with an inner cylindrical obstacle," Phys. Fluids, vol. 31,pp. 103607, 2019.
- [15] K. M. Lakshmi, P. G. Siddheshwar, D. Laroze, "Natural convection of water-copper nanoliquids confined in low-porosity cylindrical annuli," Chin. J. Phys, vol. 68, no. 4, pp. 21-136, 2020.
- [16] F. Mebarek-Oudina, "Convective heat transfer of Titania nanofuids of diferent base fuids in cylindrical annulus with discrete heat source," Heat Transf.-Asian Res, vol. 48, no. 1, pp. 135–147, 2019.
- [17] P. G. Siddheshwar, K. M. Lakshmi, "Darcy-Bénard convection of Newtonian liquids and Newtonian nanoliquids in cylindrical enclosures and cylindrical annuli," Phys. Fluids, vol. 31, pp. 084102, 2019.
- [18] F. Mebarek-Oudina, N. Keerthi Reddy, M. Sankar, "Heat source location efects on buoyant convection of nanofuids in an annulus," In Advances in Fluid Dynamics(eds Rushi Kumar, B. et al.), pp. 923–937 (Springer, 2021), 2020.
- [19] N. Keerthi Reddy, M. Sankar, "Buoyant convective transport of nanofuids in a non-uniformly heated annulus," J. Phys.: Conf. Ser, 1597, 012055, 2020.
- [20] N. Keerthi Reddy, H. A. K. Swamy, M. Sankar, "Buoyant convective fow of diferent hybrid nanoliquids in a non-uniformly heated annulus," Eur. Phys. J. Spec. Top, vol. 230, pp. 1213–1225, 2021.
- [21] K. M. Lakshmi, D. Laroze, & P. G. Siddheshwar, "A study of the natural convection of water-AA7075 nanoliquids in low-porosity cylindrical annuli using a local thermal non-equilibrium model," Heat Mass Transf, vol. 33, pp. 032018, 2021.
- [22] P.G. Siddheshwar, K.M. Lakshmi, "Darcy-Bénard convection of Newtonian liquids and Newtonian nanoliquids in cylindrical enclosures and cylindrical

annuli," Physics of Fluids, vol. 3, no. 8, pp. 084102. 2019.

- [23] K. M. Lakshmi, P. G. Siddheshwar, D. Laroze, "Natural convection of water-copper nanoliquids confined in low-porosity cylindrical annuli," Chinese Journal of Physics, vol. 68, pp. 121-136, 2020.
- [24] K. M. Lakshmi, P. G. Siddheshwar, M. S. Muddamallappa, "Study of rotating Bénard-Brinkman convection of Newtonian liquids and nanoliquids in enclosures," International Journal of Mechanical Sciences, vol. 188, pp. 105931, 2020.
- [25] M.S ankar, N.K. Reddy, Y. Do, "Conjugate buoyant convective transport of nanofluids in an enclosed annular geometry," Sci Rep, vol. 11, pp. 17122, 2021.
- [26] K. F. Sultan, "Numerical solution of heat transfer and fow of nanofuids in annulus with fns attached on the inner cylinder," J. Babylon. Univ./Eng. Sci, vol. 23, no. 2, pp. 465–484, 2015.
- [27] N. C. Roy, "Natural convection of nanofuids in a square enclosure with diferent shapes of inner geometry," Phys. Fluids, vol. 30, pp. 113605, 2018.
- [28] M.A.R. Sharif, "Laminar mixed convection in shallow inclined driven cavities with hot moving lid on top and cooled from bottom," Appl. Therm. Eng, vol. 27, pp. 1036-1042, 2007.
- [29] A. Karimipour, M. Afrand, M.M. Bazofti, "Periodic mixed convection of a nanofluid in a cavity with top lid sinusoidal motion," Int. J. Mech. Mater. Eng, vol. 1, no. 1, pp. 34-39, 2010.
- [30] E. Abu-Nada, Z. Masoud, and A. Hijazi, "Natural Convection Heat Transfer Enhancement in Horizontal Concentric Annuli Using Nanofluids," Int. Commun. Heat Mass Transfer, vol. 35, no. 5, pp. 657–665. 2008.
- [31] B.S. Kim, D.S. Lee, M.Y. Ha, H.S. Yoon, "A numerical study of natural convection in a square enclosure with a circular cylinder at different vertical locations," International Journal of Heat and Mass Transfer, vol. 51, pp. 1888–1906, 2008.
- [32] A. Ramiar, A.A. Ranjbar, S.F. Hosseinizadeh, "Effect of axial conduction and variable properties on two dimensional conjugate heat transfer of Al<sub>2</sub>O<sub>3</sub> eEG/ water mixture nanofluid in microchannel," J. Appl. Fluid Mech, vol. 5, no. 3, pp. 79-87, 2012.
- [33] K. Khanafer, K. Vafai, "A critical synthesis of thermophysical characteristics of nanofluids," Int. J. Heat Mass Transf, vol. 54, pp. 4410-4428, 2011.
- [34] H. Brinkman, "The Viscosity of Concentrated Suspensions and Solutions," J. Chem. Phys, vol. 20, no. 4, pp. 571–571, 1952.
- [35] C. Maxwell, "A Treatise on Electricity and Magnetism," Clarendon Press, Oxford, UK, (1881).

### **Creative Commons Attribution License 4.0** (Attribution 4.0 International, CC BY 4.0)

This article is published under the terms of the Creative Commons Attribution License 4.0 <u>https://creativecommons.org/licenses/by/4.0/deed.en\_US</u>