

Investigation of $(m,2)$ -methods for Solving Stiff Problems

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Abstract— A family of $(m,2)$ -methods for stiff problems solving is studied. Numerical schemes of the second and the third order are constructed. It is shown, that a maximum order of $(m,2)$ -methods is four. A -stable and L -stable numerical formulas of a maximum order are designed. An inequality for accuracy control for the L -stable methods of the fourth order is constructed. Numerical results confirming the efficiency of the constructed method are given.

Keywords— stiff system, a class of (m,k) -methods, L -stability, accuracy control of calculations.

I. INTRODUCTION

When solving the Cauchy problem for a stiff system of ordinary differential equations, Rosenbrock type methods [1] are widely used due to simple implementation and reasonably good accuracy and stability. The Rosenbrock type methods where the same Jacobi matrix is used in calculating each stage are in the most common use. It is known (see, e.g., [2]), that in this case the maximum accuracy order of the m -stage Rosenbrock method is $(m+1)$, in addition, a scheme of the maximum order can be only A -stable. If the maximum order is not required, a L -stable numerical formula of order m can be constructed. In practical calculations, as a rule, the maximum order is abandoned in favour of L -stability.

A scheme of order higher than two with the Jacobi matrix freezing can not be constructed on the basis of methods of the Rosenbrock type [3]. This limits the application of these methods to calculations with moderate accuracy or to problems of a small dimension.

In [4–5] a class of (m,k) -methods where determining a stage does not involve calculation of a right-hand part of a system of differential equations is proposed. The implementation of (m,k) -methods is as simple as that of Rosenbrock methods, however, (m,k) -schemes have advantages for accuracy and stability. In the framework of (m,k) -methods the problem of freezing the Jacobi matrix and its numerical approximation is solved more easily.

In this paper $(m,2)$ -methods for solving stiff systems where at each step the right-hand part of a system of ordinary differential equations is calculated two times are investigated. It is shown that the maximum order of accuracy of the L -stable $(m,2)$ -method is four. An inequality for accuracy control for the L -stable methods of the fourth order is constructed. Numerical results confirming the efficiency of the constructed method are given.

II. METHODS OF ROSENBRCK TYPE

We consider a Cauchy problem for a system of ordinary differential equations

$$y' = f(t, y), \quad y(t_0) = y_0, \quad t_0 \leq t \leq t_k, \quad (1)$$

where y and f are real N -dimensional vector functions, t is an independent variable, which varies over the given finite interval. It is known, that introducing an additional variable one can reduce a non-autonomous system to an autonomous form. Hence, the formulation of (1) can be considered without the loss of generality. Methods of a Rosenbrock type for problem (1) have the form

$$y_{n+1} = y_n + \sum_{i=1}^m p_i k_i, \quad (2)$$

$$D_n k_i = hf \left(y_n + \sum_{j=1}^{i-1} \beta_{ij} k_j \right),$$

$$D_n = E - ahf'_n,$$

where h is an integration step, E is an identity matrix, $f'_n = \partial f(y_n) / \partial y$ is the Jacobi matrix of a vector function $f(y)$, a , p_i , $1 \leq i \leq m$, and β_{ij} , $1 \leq i \leq m$, $1 \leq j \leq i-1$, are numerical coefficients. At present, the methods of a Rosenbrock type are treated in a wider sense [2]. Numerical formula (2) can be obtained from a class of semi-explicit methods of a Runge-Kutta type provided that only one iteration step of the Newton's method is performed when calculating each stage. In Rosenbrock methods only linear systems of algebraic equations are solved when calculating a stage, whereas in implicit or semi-explicit Runge-Kutta methods an iterative process of the Newton type is required that leads to additional problems in its implementation.

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III. THE CLASS OF (M, K) -METHODS

A class of (m, k) -methods [4] is introduced as follows. Let integer positive numbers m and k , $k \leq m$, be given. Denote a set of integer numbers i , $1 \leq i \leq m$, by M_m and the subsets of M_m of the form

$$\begin{aligned} M_k &= \{m_i \in M_m \mid 1 = m_1 < \dots < m_k \leq m\}, \\ J_i &= \{m_{j-1} \in M_m \mid j > 1, m_j \in M_k, m_j \leq i\}, \\ &2 \leq i \leq m, \end{aligned} \quad (3)$$

by M_k and J_i . Then, (m, k) -methods can be represented in the form

$$\begin{aligned} y_{n+1} &= y_n + \sum_{i=1}^m p_i k_i, \quad D_n = E - ahf'_n, \\ D_n k_i &= hf \left(y_n + \sum_{j=1}^{i-1} \beta_{ij} k_j \right) + \sum_{j \in J_i} \alpha_{ij} k_j, \quad i \in M_k, \\ D_n k_i &= k_{i-1} + \sum_{j \in J_i} \alpha_{ij} k_j, \quad i \in M_m \setminus M_k. \end{aligned} \quad (4)$$

The set J_i , $2 \leq i \leq m$, serves to eliminate the redundant coefficients α_{ij} by means of which we can not make effect on accuracy and stability of (4) and which can be expressed linearly in terms of other coefficients. In traditional one-step methods, to determine computational work per integration step a single constant m being the number of stages is sufficient because in these methods each stage is accompanied by obligatory calculation of the right-hand part of (1). In methods (4) there are two types of stages. For some stages it is necessary to calculate the right-hand part, but for other ones it is not required. As a result, to determine computational work per step in (4) two constants m and k are needed. The expense of a step is as follows: the Jacobi matrix is calculated once and the decomposition of a matrix D_n is performed once, function f is calculated k times, backward Gauss is performed m times. For $k = m$ and $\alpha_{ij} = 0$ numerical schemes (4) coincide with methods (2) of a Rosenbrock type. In other cases these methods differ and methods (4) have advantages over methods (2).

IV. THE MAXIMAL ORDER OF ACCURACY OF $(M, 2)$ -METHODS

We consider $(m, 2)$ -methods of the following form

$$\begin{aligned} y_{n+1} &= y_n + \sum_{i=1}^m p_i k_i, \\ D_n &= E - ahf'_n, \\ D_n k_1 &= hf(y_n), \quad D_n k_i = k_{i-1}, \quad 2 \leq i \leq s_1 - 1, \\ D_n k_{s_1} &= hf \left(y_n + \sum_{j=1}^{s_1-1} \beta_{s_1, j} k_j \right) + \alpha_{s_1, s_1-1} k_{s_1-1}, \\ D_n k_i &= k_{i-1} + \alpha_{i, s_1-1} k_{s_1-1}, \quad s_1 + 1 \leq i \leq m, \end{aligned} \quad (5)$$

where s_1 and m , $s_1 \leq m$, are arbitrary integer constants. It is easy to see, that formulas (5) describe all kinds of $(m, 2)$ -methods (4).

Theorem. For any number m of stages and for any set M_2 it is impossible to construct an $(m, 2)$ -method of accuracy of order higher than four.

Without the loss of generality we give a proof for scalar problem (1), which exact solution $y(t_{n+1})$ can be written in the form

$$\begin{aligned} y(t_{n+1}) &= y(t_n) + hf + \frac{1}{2} h^2 ff \\ &+ \frac{1}{6} h^3 [f'^2 f + f''f^2] \\ &+ \frac{1}{24} h^4 [f'^3 f + 4ff''f^2 + f''f^3] \\ &+ \frac{1}{120} h^5 [f'^4 f + 4ff''f^3 + 5f'^2 f''f^2 + f''^2 f^3 + f^{IV} f^4] \\ &+ O(h^6), \end{aligned} \quad (6)$$

where the elementary differentials are calculated on an exact solution $y(t_n)$. We consider $(m, 2)$ -methods (5). Taking into account

$$\begin{aligned} D_n^{-1} &= E + ahf'_n + a^2 h^2 f_n'^2 + a^3 h^3 f_n'^3 \\ &+ a^4 h^4 f_n'^4 + O(h^5), \end{aligned} \quad (7)$$

we observe that the second calculation of a function $f(y)$ is performed at the point

$$y_{n,c} = y_n + \sum_{j=1}^{s_1-1} \beta_{s_1, j} k_j = y_n + \sum_{i=1}^5 c_i h^i f_n'^{(i-1)} f_n + O(h^6),$$

where c_i , $1 \leq i \leq 4$, are determined in terms of the coefficients of scheme (5) and the elementary differentials are calculated on an approximate solution y_n . Taking into consideration (6) and (7), to prove the theorem it is sufficient to show that the Taylor expansion of the function $f(y_{n,c})$ in terms of powers of h does not involve the term $h^5 f_n''^2 f_n^3$. The Taylor expansion of $f(y_{n,c})$ in a vicinity of the point y_n up to terms of order h^5 inclusive has the form

$$\begin{aligned} f(y_{n,c}) &= hf_n + c_1 h^2 f_n' f_n + h^3 \left[c_2 f_n'^2 f_n + \frac{1}{2} f_n'' f_n^2 \right] \\ &+ h^4 \left[c_3 f_n'^3 f_n + c_1 c_2 f_n' f_n'' f_n^2 + \frac{1}{6} c_1^3 f_n''^3 f_n^3 \right] \\ &+ h^5 \left[c_4 f_n'^4 f_n + c_1 c_3 f_n'^3 f_n'' f_n^2 + \frac{1}{2} c_1^2 c_2 f_n' f_n'' f_n^3 + \frac{1}{24} c_1^4 f_n^{IV} f_n^4 \right] \\ &+ O(h^6), \end{aligned}$$

which completes the proof.

V. AN A-STABLE (M,2)-METHOD OF ORDER FOUR

Let $k = 2$. Take $M_k = \{1, 2\}$, then $J_i = \{1\}$, $2 \leq i \leq m$. As a result, we obtain the numerical schemes of the form

$$\begin{aligned} y_{n+1} &= y_n + \sum_{i=1}^m p_i k_i, \\ D_n k_1 &= hf(y_n), \\ D_n k_2 &= hf(y_n + \beta_{21} k_1) + \alpha_{21} k_1, \\ D_n k_i &= k_{i-1} + \alpha_{i1} k_1, \quad 3 \leq i \leq m. \end{aligned} \quad (8)$$

Case 1. Let $m = 2$. Then the coefficients

$$\begin{aligned} a &= \frac{6 + \sqrt{12}}{12}, \quad p_1 = \frac{76a - 3}{54a}, \quad p_2 = \frac{16}{27}, \\ \beta_{21} &= \frac{3}{4}, \quad \alpha_{21} = \frac{3 - 54a}{32a} \end{aligned}$$

provide the third order of the A -stable scheme (8). Usually the maximal order is abandoned in favour of L -stability. For $m = 2$ the coefficients of the second order L -stable scheme have the form

$$a = \beta_{21} = 1 - \frac{\sqrt{2}}{2}, \quad p_1 = a, \quad p_2 = 1 - a, \quad \alpha_{21} = 0.$$

Case 2. Let $m = 3$. Introduce the notations

$$\begin{aligned} c_1 &= 1 - 4a + 2a^2, \quad c_2 = \beta_{21} + a\alpha_{21}, \\ c_3 &= 6ac_2\beta_{21}^2, \quad c_4 = 2c_2(\beta_{21} - a), \\ c_5 &= 3c_1\beta_{21}^2(\beta_{21} - a). \end{aligned}$$

Then the coefficients

$$\begin{aligned} p_1 &= \frac{ac_3 + c_4 - c_5}{c_3}, \\ p_2 &= \frac{2c_2 - 3\beta_{21}^2 c_1}{6\beta_{21}^2 c_2}, \quad p_3 = \frac{c_1}{2c_2}, \\ \alpha_{31} &= \frac{2c_2[(1-a)c_3 - c_4 + c_5 - 2ac_2(1 + \alpha_{21})]}{c_2 c_3}. \end{aligned}$$

provide third accuracy order of the L -stable scheme (8). Here, α_{21} and β_{21} are constant coefficients and the value of a is determined from the equation

$$6a^3 - 18a^2 + 9a - 1 = 0.$$

It is easy to verify that for any value of m it is impossible to construct an L -stable scheme (8) of the fourth order. To do this, it is sufficient to write conditions of the fourth accuracy order consistency and L -stability. The simplest study of the obtained nonlinear system of algebraic equations shows its incompatibility. However, provided L -stability is not required, the method (8) of fourth accuracy order can be constructed.

Case 3. Let $m = 4$. Introduce the notations

$$\begin{aligned} c_1 &= \frac{60 - 81c_7}{81}, \quad c_2 = \frac{18 - 324c_7}{81}, \\ c_3 &= \frac{405c_7 - 64}{81}, \quad c_4 = \frac{19 - 162c_7}{81}, \end{aligned}$$

$$c_5 = \frac{64 - 81c_7}{81}, \quad c_6 = \frac{162c_7 - 16}{81},$$

where $c_7 \neq 0$ is a constant coefficient. Then the coefficients of the fourth order method have the form

$$\begin{aligned} a &= \frac{3}{8}, \quad \beta_{21} = \frac{3}{4}, \quad p_1 = c_1, \quad p_2 = c_5, \quad p_3 = c_6, \\ p_4 &= c_7, \quad \alpha_{21} = \frac{c_4}{c_7}, \quad \alpha_{31} = \frac{c_3 c_7 - c_4 c_6}{c_7^2}, \\ \alpha_{41} &= \frac{c_2 c_7^2 - c_4 c_5 c_7 - c_3 c_6 c_7 + c_4 c_6^2}{c_7^2}. \end{aligned}$$

VI. AN L-STABLE (M,2)-METHOD OF ORDER FOUR

Let $k = 2$. Take $M_k = \{1, 3\}$, then $J_i = \{2\}$, $3 \leq i \leq 4$. As a result, we get a scheme of the form

$$\begin{aligned} y_{n+1} &= y_n + \sum_{i=1}^4 p_i k_i, \\ D_n k_1 &= hf(y_n), \quad D_n k_2 = k_1, \\ D_n k_3 &= hf(y_n + \beta_{31} k_1 + \beta_{32} k_2) + \alpha_{32} k_2, \\ D_n k_4 &= k_3 + \alpha_{42} k_2. \end{aligned} \quad (9)$$

Construct an L -stable method of order four. To do this, substitute the expansion of the stages k_i , $1 \leq i \leq 4$, in the Taylor series in the first formulae of (9). Assuming that $y_n = y(t_n)$ and comparing the obtained presentation of the approximate solution y_{n+1} with the Taylor expansion of the exact solution $y(t_{n+1})$, write the conditions of the fourth accuracy order

$$\begin{aligned} p_1 + p_2 + (1 + \alpha_{32}) p_3 + (1 + \alpha_{32} + \alpha_{42}) p_4 &= 1, \\ ap_1 + 2ap_2 + (a + \beta_{31} + \beta_{32} + 3a\alpha_{32}) p_3 &+ \\ + (2a + \beta_{31} + \beta_{32} + 4a\alpha_{32} + 3a\alpha_{42}) p_4 &= \frac{1}{2}, \\ a^2 p_1 + 3a^2 p_2 + (a^2 + 2a\beta_{31} + 3a\beta_{32} + 6a^2\alpha_{32}) p_3 &+ \\ + (3a^2 + 3a\beta_{31} + 4a\beta_{32} + 10a^2\alpha_{32} + 6a^2\alpha_{42}) p_4 &= \frac{1}{6}, \\ a^3 p_1 + 4a^3 p_2 + (a^3 + 3a^2\beta_{31} + 6a^2\beta_{32} + 10a^3\alpha_{32}) p_3 &+ \\ + (4a^3 + 6a^2\beta_{31} + 10a^2\beta_{32} + 20a^3\alpha_{32} + 10a^3\alpha_{42}) p_4 &= \frac{1}{24}, \\ (\beta_{31} + \beta_{32})^2 (p_3 + p_4) &= \frac{1}{3}, \\ a(\beta_{31} + \beta_{32})(\beta_{31} + 2\beta_{32})(p_3 + p_4) &= \frac{1}{8}, \\ a(\beta_{31} + \beta_{32})^2 \left(\frac{1}{2} p_3 + p_4 \right) &= \frac{1}{24}, \\ (\beta_{31} + \beta_{32})^3 (p_3 + p_4) &= \frac{1}{4}. \end{aligned}$$

Applying the method (9) to the solution of the test problem

$$y' = \lambda y, \quad y(0) = y_0, \quad t \geq 0,$$

obtain the condition of L -stability of the numerical formulae (9) of the form

$$a(a - p_1) + (\beta_{31} - a)p_3 = 0.$$

Investigating the consistency of this relation and the order conditions, write

$$p_1 = \frac{76a^2 - 29a + 3}{27a^2}, \quad p_2 = \frac{146a^2 + 89a - 12}{27a^2},$$

$$p_3 = \frac{32a - 4}{27a}, \quad p_4 = \frac{4 - 16a}{27a}, \quad \beta_{31} = \frac{48a - 9}{32a},$$

$$\beta_{32} = \frac{9 - 24a}{32a}, \quad \alpha_{32} = \frac{-54a^2 + 57a - 12}{8a - 32a^2},$$

$$\alpha_{42} = \frac{-864a^3 + 828a^2 - 288a + 36}{a(4 - 16a)^2},$$

where a is determined from the necessary condition of L -stability

$$24a^4 - 96a^3 + 72a^2 - 16a + 1 = 0.$$

The given equation have four real roots

$$a_1 = 0.10643879214266, \quad a_2 = 0.22042841025921,$$

$$a_3 = 0.57281606248213, \quad a_4 = 3.10031673511599.$$

In calculations it is wise to take $a = 0.57281606248213$, and the coefficients of the scheme (9) have the form

$$p_1 = 1.27836939012447, \quad p_2 = -1.00738680980438, \\ p_3 = 0.92655391093950, \quad p_4 = -0.33396131834691, \quad (10)$$

$$\beta_{31} = 1.00900469029922, \quad \beta_{32} = -0.25900469029921,$$

$$\alpha_{32} = -0.49552206416578, \quad \alpha_{42} = -1.28777648233922.$$

VII. ACCURACY CONTROL

The third accuracy order method of the form

$$y_{n+1,1} = y_n + b_1k_1 + b_2k_2 + b_3k_3 + b_4k_4$$

can be applied for accuracy control of numerical formula (9) of the fourth order. Here, $D_n k_5 = k_4$, and k_i , $1 \leq i \leq 3$, are determined in (9). It is easy to see, that the third order requirements have the form

$$b_1 + b_2 + (1 + \alpha_{32})b_3 + (1 + \alpha_{32} + \alpha_{42})b_4 = 1,$$

$$ab_1 + 2ab_2 + (a + \beta_{31} + \beta_{32} + 3a\alpha_{32})b_3$$

$$+ (3a + \beta_{31} + \beta_{32} + 5a\alpha_{32} + 4a\alpha_{42})b_4 = \frac{1}{2},$$

$$a^2b_1 + 3a^2b_2 + (a^2 + 2a\beta_{31} + 3a\beta_{32} + 6a^2\alpha_{32})b_3$$

$$+ (6a^2 + 4a\beta_{31} + 5a\beta_{32} + 15a^2\alpha_{32} + 10a^2\alpha_{42})b_4 = \frac{1}{6},$$

$$(\beta_{31} + \beta_{32})^2 (b_3 + b_4) = \frac{1}{3},$$

where the coefficients a , β_{31} , β_{32} , α_{32} and α_{42} are set in (10). The given system is linear with respect to b_i , $1 \leq i \leq 4$.

As a result, we have

$$b_1 = 1.203100567018353,$$

$$b_2 = -6.552116304144386 \cdot 10^{-1},$$

$$b_3 = 7.115271884598151 \cdot 10^{-1},$$

$$b_4 = -1.189345958672225 \cdot 10^{-1}.$$

Now, we can calculate an error estimation ε_n by the formula

$$\varepsilon_n = \|y_{n+1} - y_{n+1,1}\|,$$

and verify the inequality $\varepsilon_n \leq \varepsilon$ on a choice of an integration step size, where $\|\cdot\|$ is some norm in R^N , ε is accuracy of calculations.

VIII. THE RESULTS OF CALCULATIONS

The following example describes a system of two differential equations in partial derivatives with initial and boundary conditions. The Akzo Nobel Central Research laboratory has formulated a research problem of the penetration of a labeled tracer antibody in the tumor tissue of a living organism [7]. A system of one-dimensional reaction-diffusion equations

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} - kuv, \quad \frac{\partial v}{\partial t} = -kuv, \quad (11)$$

which arise from the chemical reaction $A + B \xrightarrow{k} C$, where A is an antibody with a radioactive label, which reacts with substrate B , the tissue, affected by the tumor, and k is the reaction rate constant. Concentrations of A and B are denoted by u and v , respectively. In deriving equations (11), we assumed that the kinetics of the reaction is described by the mass action law, and reagent A is mobile, whereas reagent B is fixed. We study a semi-infinite plate, inside which substrate B is uniformly distributed. Reagent A , when applied to the surface of the plate, begins to penetrate into it. To simulate the penetration of equation (11), in the strip, the following conditions

$$S_T = \{(x, t) : 0 < x < \infty, 0 < t < T\},$$

with the initial conditions

$$u(x, 0) = 0, \quad v(x, 0) = v_0, \quad x > 0$$

and the boundary conditions

$$u(0, t) = \varphi(t), \quad 0 < t < T,$$

where v_0 is a constant – are considered. For the numerical solution, the variable x is converted so that the semi-infinite plate is transformed into the final one. This transformation provides a special family of transformations of Mobius

$$\zeta = \frac{x}{x+c}, \quad c > 0.$$

Each transformation of this family transforms S_T into rectangle

$$\{(\zeta, t) : 0 < \zeta < 1, 0 < t < T\}.$$

Using ζ , problem (11) is rewritten as

$$\frac{\partial u}{\partial t} = \frac{(\zeta - 1)^4}{c^2} \frac{\partial^2 u}{\partial \zeta^2} + \frac{2(\zeta - 1)^3}{c^2} \frac{\partial u}{\partial \zeta} - kuv,$$

$$\frac{\partial v}{\partial t} = -kuv,$$

with initial conditions

$$u(\zeta, 0) = 0, v(\zeta, 0) = v_0, \zeta > 0$$

and the boundary conditions

$$u(0, t) = \varphi(t), \frac{\partial u(1, t)}{\partial \zeta} = 0, 0 < t < T,$$

The last boundary condition is obtained from the ratio $\partial u(\infty, t) / \partial x = 0$. Discretization of the derivatives in the spatial variables using the method of straight lines leads to the Cauchy problem for a system of ordinary differential equations. For discretization, we used a uniform grid

$$\{\zeta_j\}, \zeta_j = j \cdot \Delta \zeta, \Delta \zeta = \frac{1}{N}, 1 \leq j \leq N.$$

Variables u_j and v_j denote the approximations $u(\zeta_j, t)$ and $v(\zeta_j, t)$, respectively. It is clear that u_j and v_j are functions of t . The discretizations of the first and second order derivatives on the spatial variable, respectively, have the form

$$\frac{\partial u_j}{\partial \zeta} = \frac{u_{j+1} - u_{j-1}}{2\Delta \zeta},$$

$$\frac{\partial^2 u_j}{\partial \zeta^2} = \frac{u_{j-1} - 2u_j + u_{j+1}}{(\Delta \zeta)^2},$$

$$1 \leq j \leq N.$$

The values of u_0 and u_{N+1} , obtained from the boundary conditions, have the form $u_0 = \varphi(t)$ and $u_{N+1} = u_N$. For

$$y = (u_1, v_1, u_2, v_2, \dots, u_N, v_N)^T$$

and $T = 20$, this semidiscretization problem can be written as

$$\frac{dy}{dt} = f(t, y), y(0) = g, \quad (12)$$

$$y \in R^{2N}, 0 \leq t \leq 20,$$

where N is set by a user parameter. The function f is determined by formulas

$$f_{2j-1} = \alpha_j \frac{y_{2j+1} - y_{2j-3}}{2\Delta \zeta} + \beta_j \frac{y_{2j-3} - 2y_{2j-1} + y_{2j+1}}{(\Delta \zeta)^2} - ky_{2j-1}y_{2j},$$

$$f_{2j} = -ky_{2j}y_{2j-1},$$

where

$$\alpha_j = \frac{2}{c^2}(j\Delta \zeta - 1)^3, \beta_j = \frac{1}{c^2}(j\Delta \zeta - 1)^4, 1 \leq j \leq N,$$

$$\Delta \zeta = \frac{1}{N}, y_{-1}(t) = \varphi(t), y_{2, N+1} = y_{2, N-1},$$

$$g \in R^{2N}, g = (0, v_0, 0, v_0, \dots, 0, v_0)^T.$$

Function $\varphi(t) = 2$ for $t \in (0, 5]$ and $\varphi(t) = 0$ for $t \in (5, 20]$, that is, φ has a first-order discontinuity at the point $t = 5$. According to [7], the appropriate values for the parameters k , v_0 and c are $k = 100$, $v_0 = 1$ and $c = 4$. It also shows the results of calculations with high accuracy and extended bit grating.

The calculations were performed with accuracy $\varepsilon = 10^{-4}$. The following calculations were performed with the numerical Jacobi matrix for $N = 200$, i.e., system (12) consists of 400 equations. The task of finding the gap of function $\varphi(t)$ at $t = 5$ was assigned to the control algorithm of the step. In this paper i_f and i_j denote, respectively, the total numbers of calculations of the right-hand part and the Jacobi matrix decompositions of the problem (12), which allow to estimate the efficiency of the integration algorithm objectively. The solution of this problem by algorithm was calculated with costs $i_f = 76\,717$ and $i_j = 95$.

IX. CONCLUSION

From the above considerations we can draw the following conclusions.

Firstly, stability of (m, k) -methods depends on a choice of sets (3), i.e., on a way of implementation of numerical schemes. This follows from a comparison of formulas (8) and (9).

Secondly, a $(4, 2)$ -method, which is competitive in accuracy with the implicit method of a Runge-Kutta type with two calculations of the right-hand part of (1) can be constructed. In linear analysis of stability this method is also as good as the implicit method of a Runge-Kutta type. In addition, in (9) a way of implementation is included, i.e., one can estimate computational work per integration step before calculations. Computational work of the implicit methods of a Runge-Kutta type depends highly on a way of implementation. The use of two-stage scheme does not imply that at each step the right-hand part of (1) is calculated twice. Therefore, for some problems the $(4, 2)$ -method is preferred over implicit numerical formulas of the Runge-Kutta type.

Thirdly, when a function $f(y)$ of problem (1) is calculated twice, an L -stable $(4, 2)$ -method of the fourth accuracy order can be constructed whereas corresponding L -stable method (2) of the Rosenbrock type can be of order two only. For sufficiently high accuracy of calculations and large dimension of problem (1), decomposition of the Jacobi matrix is responsible, in fact, for total computational work, whereas impact of backward Gauss is unessential. In this case $(4, 2)$ -method is more efficient.

Notice, that in the framework of (m, k) -methods the problem of freezing the Jacobi matrix can be solved easily [6].

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