

# Researches concerning the dynamical flow and the reaction products engendering into a burning installation as variation of the methane gas injected into furnace

Mihai D.L. Țălu and Ștefan D.L. Țălu

**Abstract**—In the paper there are described the analysis of the dynamic flow and the reaction products engendering as consequence of the methane variation injected through a burner into the furnace. The burning installation has two identical burners. This research work is fulfilled using the FLUENT programme.

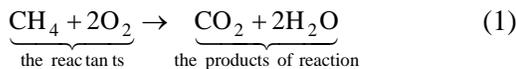
**Keywords**—Burner, methane combustion, Finite Element Method, computational simulation.

## I. INTRODUCTION

THE burning installation on methane gas have one or more injectors [1].

To establish the field of gas-dynamical flow of methane into a burner it is necessary to solve the transport equations of reactants and products of reaction [2]-[7].

The mechanism of irreversible reaction with complete conversion for methane oxidation reaction is:



The mathematical model which describes the tensorial form of flow consists by [6]:

- the Reynolds flow turbulent equation (the flame into the industrial installations is designed to be in a turbulent flow regime):

$$\begin{aligned} \frac{\partial}{\partial t}(\bar{\rho}\bar{v}_i) + \frac{\partial}{\partial x_j}(\bar{\rho}\bar{v}_j\bar{v}_i) + \frac{\partial \bar{p}}{\partial x_j} - \bar{\rho}\bar{f}_i + \\ + \frac{\partial}{\partial x_j} \left\{ \mu_{\text{ef}} \left[ \left( \frac{\partial \bar{v}_i}{\partial x_j} + \frac{\partial \bar{v}_j}{\partial x_i} \right) - \frac{2}{3} \frac{\partial \bar{v}_i}{\partial x_j} \delta_{ij} \right] \right\} = 0 \end{aligned} \quad (2)$$

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- the kinetic energy of turbulence:

$$\frac{\partial}{\partial t}(\bar{\rho}k) + \frac{\partial}{\partial x_j} \left( \bar{\rho}\bar{v}_j k - \frac{\mu_{\text{ef}}}{\sigma_k} \frac{\partial k}{\partial x_j} \right) + c\bar{p}\epsilon - \quad (3)$$

$$- \frac{\partial v_i}{\partial x_j} \left\{ \mu_{\text{ef}} \left[ \left( \frac{\partial \bar{v}_i}{\partial x_j} + \frac{\partial \bar{v}_j}{\partial x_i} \right) - \frac{2}{3} k \delta_{ij} \right] \right\} = 0$$

- the mass continuity equation:

$$\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial}{\partial x_j}(\bar{\rho}\bar{v}_j) = 0 \quad (4)$$

- the transport equations for reactants CH<sub>4</sub> and O<sub>2</sub>:

$$\frac{\partial}{\partial t}(\bar{\rho} \cdot m_{\text{CH}_4}) + \nabla \cdot (\bar{\rho} \cdot \bar{v} \cdot m_{\text{CH}_4}) - \nabla \cdot (\bar{\rho} \cdot D m_{\text{CH}_4}) = S_{\text{CH}_4} \quad (5)$$

$$\frac{\partial}{\partial t}(\bar{\rho} \cdot m_{\text{O}_2}) + \nabla \cdot (\bar{\rho} \cdot \bar{v} \cdot m_{\text{O}_2}) - \nabla \cdot (\bar{\rho} \cdot D m_{\text{O}_2}) = S_{\text{O}_2} \quad (6)$$

- the transport equations for products of reaction CO<sub>2</sub>, H<sub>2</sub>O, N<sub>2</sub>:

$$\frac{\partial}{\partial t}(\bar{\rho} \cdot m_{\text{CO}_2}) + \nabla \cdot (\bar{\rho} \cdot \bar{v} \cdot m_{\text{CO}_2}) - \nabla \cdot (\bar{\rho} \cdot D m_{\text{CO}_2}) = S_{\text{CO}_2} \quad (7)$$

$$\frac{\partial}{\partial t}(\bar{\rho} \cdot m_{\text{H}_2\text{O}}) + \nabla \cdot (\bar{\rho} \cdot \bar{v} \cdot m_{\text{H}_2\text{O}}) - \nabla \cdot (\bar{\rho} \cdot D m_{\text{H}_2\text{O}}) = S_{\text{H}_2\text{O}} \quad (8)$$

$$\frac{\partial}{\partial t}(\bar{\rho} \cdot m_{\text{N}_2}) + \nabla \cdot (\bar{\rho} \cdot \bar{v} \cdot m_{\text{N}_2}) - \nabla \cdot (\bar{\rho} \cdot D m_{\text{N}_2}) = S_{\text{N}_2} \quad (9)$$

This paper presents the partial results of the 3D distribution fields concerning the dynamic flow fields of pressure, density, temperature, velocity and the distribution mass fraction for the products of reaction: CH<sub>4</sub>, CO<sub>2</sub>, O<sub>2</sub>, H<sub>2</sub>O and N<sub>2</sub> emitted into a burning chamber with two identical methane burner, Fig. 2, as result of the heat density variation of one of them (Br1).

The main constructive sizes of burning installation are given in Table I [8].

Table I. The main constructive sizes of burning installation.

r <sub>1</sub>	R <sub>1</sub>	r <sub>2</sub>	R <sub>2</sub>	l <sub>c</sub>	L	l <sub>1</sub> = l <sub>2</sub> = 2.2 (2R)	α
[mm]							[°]
3	10	12	30	18	4500	110	7.5

The burner's scheme and their 3D assembling on

installation are shown in Fig. 1 and Fig. 2.

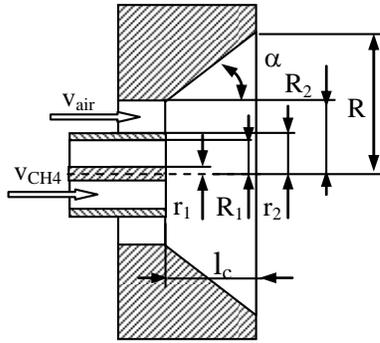


Fig. 1. The burner's scheme.

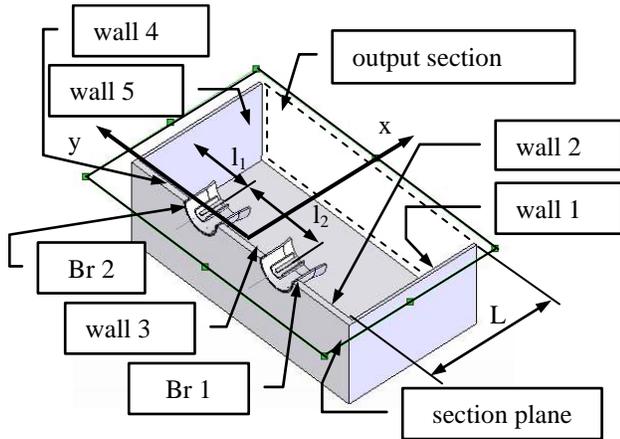


Fig. 2. The burning installation.

The coefficient of air excess is  $\lambda = 1.1$  and the temperatures  $T_{CH_4} = T_{air} = 300$  [K]; the output of burner is connected to atmosphere. The flow volumes of feeding the burners are given in Table II.

II. THE ANALYSIS WITH THE F.E.M.

The F.E.M. analysis and the modeling of the burning chamber is made using the Fluent 6.2.16 and Gambit 2.2.30 programmes. The partial results of analysis concerning the 3D fields of the dynamic flow in the plane with maximum heat density which passing through the burners axes, Fig. 2 and into 3D domain of flow are presented.

The changing of the heat density of burner Br1 (through variations of  $CH_4$  which is injected into furnace) demanded a separately study of flow fields in the interior domain of flow and on the frontiers: walls of burning chamber and on the output section, Fig. 2.

The study cases are given in Table II [8].

Table II. The study cases.

	Case	$Q_{CH_4}$	$Q_{air}$		Case	$Q_{CH_4}$	$Q_{air}$
		[m <sup>3</sup> /h]				[m <sup>3</sup> /h]	
Br1	C1	0	0		C7	9.26	103.98
	C2	4.11	47.55		C8	10.29	118.87
	C3	5.14	59.43		C9	11.32	130.75
	C4	6.17	71.32		C10	12.35	142.65
	C5	7.20	83.21	Br2	C6	8.23	95.10
	C6	8.23	95.10				

The numerical results of simulation into volume domain of flow are given in Table III and for the temperature on the frontiers in Table IV.

Table III. The numerical results of simulation into volume domain of flow.

Physical sizes	Integral volume average of sizes			
	Case			
	C1	C2	C3	C4
p[Pa]	0.0074	1.148	1.57623	2.00155
$\rho$ [kg/m <sup>3</sup> ]	0.17541	0.1809	0.18192	0.18283
T[K]	1908.39	1865.8	1857.05	1849.03
v[m/s]	7.26	10.76	11.6066	12.4443
CH <sub>4</sub>	0.0696	0.0789	0.08025	0.08061
CO <sub>2</sub>	0.13211	0.13056	0.12970	0.12889
O <sub>2</sub>	0.01079	0.01084	0.01205	0.01321
N <sub>2</sub>	0.679	0.67195	0.67180	0.67175
H <sub>2</sub> O	0.10815	0.10689	0.10618	0.10552
	C5	C6	C7	C8
p[Pa]	2.39350	2.78311	3.12201	3.44071
$\rho$ [kg/m <sup>3</sup> ]	0.18343	0.18370	0.18371	0.18350
T[K]	1843.75	1841.62	1842.11	1844.48
v[m/s]	13.2943	14.1631	15.0500	15.9558
CH <sub>4</sub>	0.08083	0.08087	0.08076	0.08058
CO <sub>2</sub>	0.12835	0.12809	0.12806	0.12819
O <sub>2</sub>	0.01399	0.01438	0.01446	0.01429
N <sub>2</sub>	0.67172	0.67176	0.67186	0.67196
H <sub>2</sub> O	0.10508	0.10487	0.10484	0.10495
	C9	C10		
p[Pa]	3.69408	3.93328		
$\rho$ [kg/m <sup>3</sup> ]	0.18317	0.18281		
T[K]	1847.97	1851.74		
v[m/s]	16.8730	17.7925		
CH <sub>4</sub>	0.08033	0.08001		
CO <sub>2</sub>	0.12843	0.12869		
O <sub>2</sub>	0.01398	0.01366		
N <sub>2</sub>	0.67209	0.67226		
H <sub>2</sub> O	0.10515	0.10536		

Considering as reference the values from case C6, in Fig. 3 the abscissae represents the percentage of effective methane

flow volume injected in burner Br1 ( $Q_{1ef}$ ) and on the ordinate the percentage of the variation for the physical sizes from cases C1 to C10 [9].

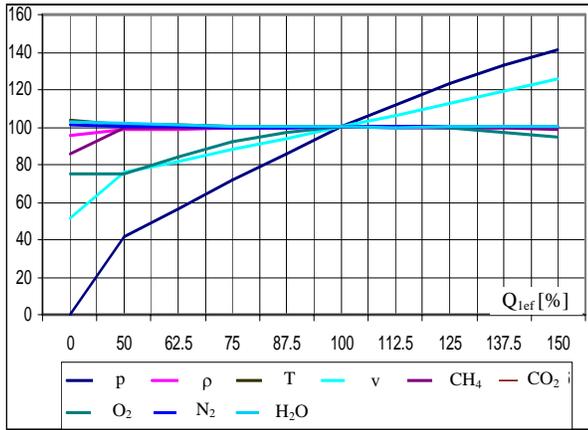


Fig. 3. The percentage of the variation for the physical sizes from cases C1 to C10 in function of  $Q_{1ef}$ .

Table IV. The numerical results of simulation for the temperature on the frontiers.

Case		T[K]	
		Medium value	Maximum value
C1	wall 1	2150.50	2292.60
	wall 2	1713.20	1738.80
	wall 3	1414.80	1535.00
	wall 4	1367.65	1445.00
	wall 5	2093.85	2308.60
	output section	1921.25	2245.80
C2	wall 1	2151.15	2303.90
	wall 2	1340.80	1457.00
	wall 3	878.53	1010.90
	wall 4	1498.30	1589.40
	wall 5	2198.75	2313.95
	output section	1912	2191.56
C3	wall 1	2166.00	2305.30
	wall 2	1324.70	1437.20
	wall 3	816.44	929.70
	wall 4	1521.16	1614.85
	wall 5	2204.36	2314.74
	output section	1913.67	2173.87
C4	wall 1	2168.31	2305.50
	wall 2	1311.60	1421.30
	wall 3	780.76	882.86
	wall 4	1540.16	1636.16
	wall 5	2205.52	2315.30
	output section	1916.54	2186.15
C5	wall 1	2162.02	2305.75
	wall 2	1300.47	1407.85
	wall 3	761.42	857.70
	wall 4	1557.00	1655
	wall 5	2204.01	2315.72
	output section	1916.85	2201.27

C6	wall 1	2150.90	2305.60
	wall 2	1289.52	1394.70
	wall 3	755.84	850.02
	wall 4	1571.11	1671.44
	wall 5	2150.90	2305.60
	output section	1918.73	2209.07
C7	wall 1	2136.46	2305.36
	wall 2	1279.59	1382.96
	wall 3	759.16	853.98
	wall 4	1584.01	1685.88
	wall 5	2196.36	2315.81
	output section	1919.03	2213.74
C8	wall 1	2120.97	2305
	wall 2	1270.75	1372.6
	wall 3	768.70	866
	wall 4	1596.10	1698.9
	wall 5	2191.04	2315.35
	output section	1920.81	2230.10
C9	wall 1	2105.56	2304.60
	wall 2	1262.84	1363.40
	wall 3	784.18	885.70
	wall 4	1606.33	1710.50
	wall 5	2185.33	2314.50
	output section	1920.93	2242.75
C10	wall 1	2090.40	2304.15
	wall 2	1255.65	1355.05
	wall 3	802.84	909.70
	wall 4	1615.76	1721.25
	wall 5	2179.65	2312.90
	output section	1921.27	2251.60

From Fig. 4 to Fig. 11 as exemplification are given the distributions fields of the dynamical flow for sizes from Table III corresponding to the case C2.

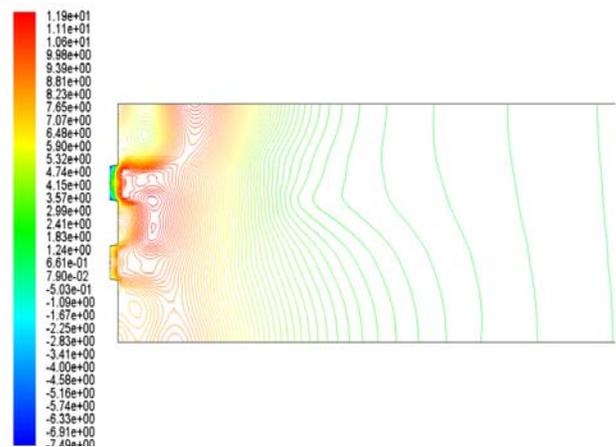


Fig. 4. The pressure.

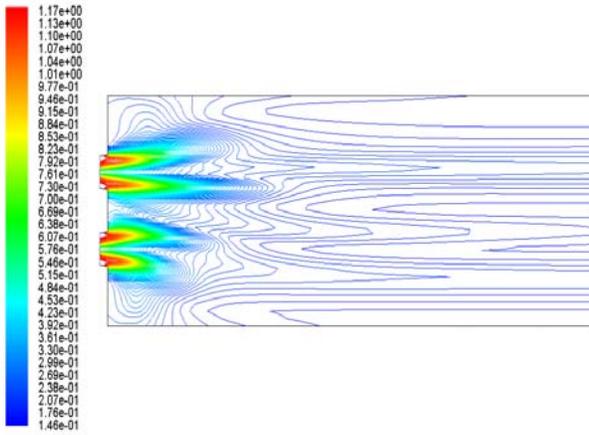


Fig. 5. The density.

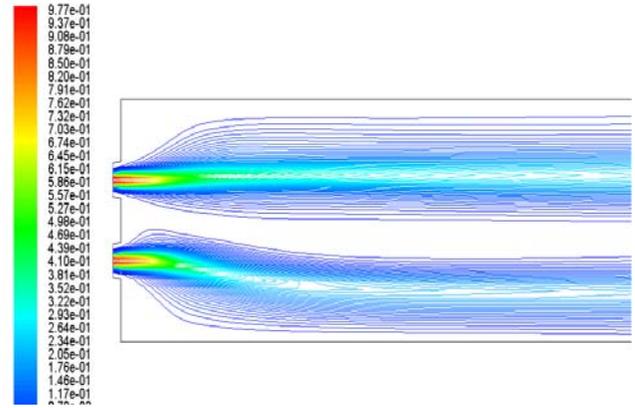


Fig. 8. The mass fraction of CH<sub>4</sub>.

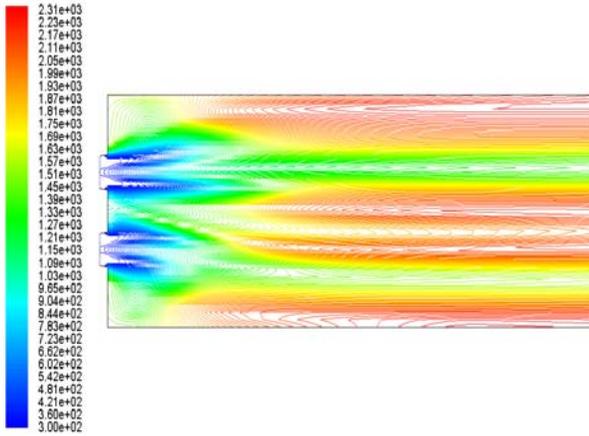


Fig. 6. The temperature.

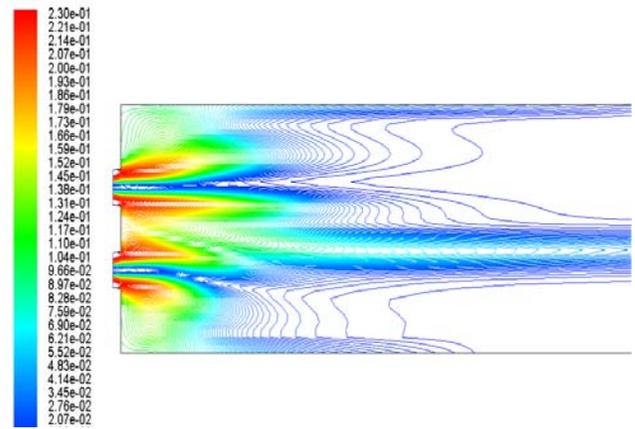


Fig. 9. The mass fraction of O<sub>2</sub>.

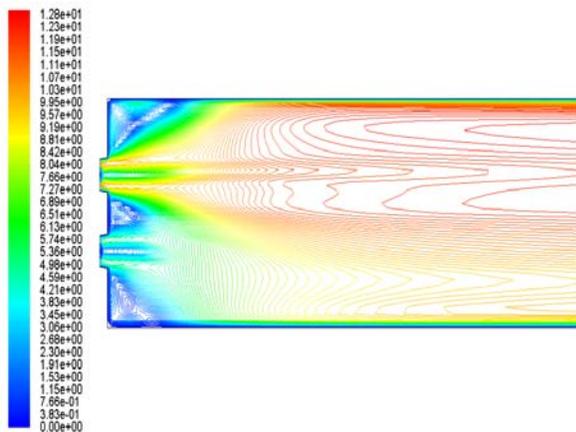


Fig. 7. The velocity.

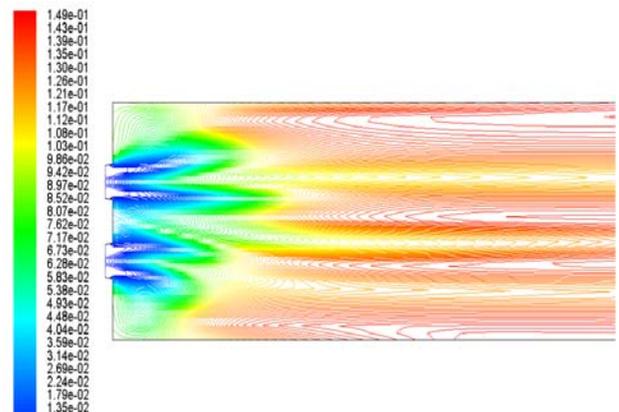


Fig. 10. The mass fraction of CO<sub>2</sub>.

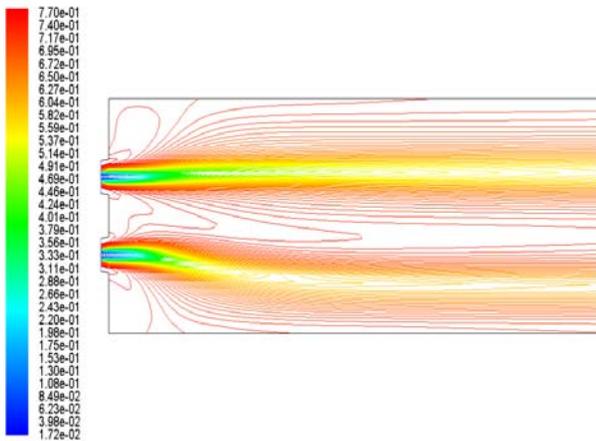


Fig. 11. The mass fraction of N<sub>2</sub>.

The temperature fields in section plane with maximum heat density for cases C1, C4, C8 and C10 are shown in Fig. 12 to Fig. 15. The medium value and the maximum value of temperature on domain frontiers are given in Table IV.

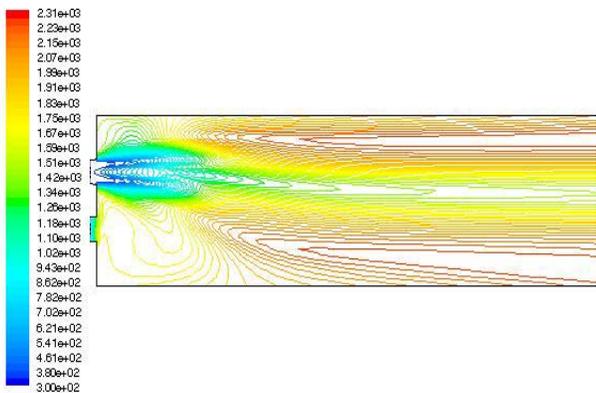


Fig. 12. The temperature fields in section plane with maximum heat density for case C1.

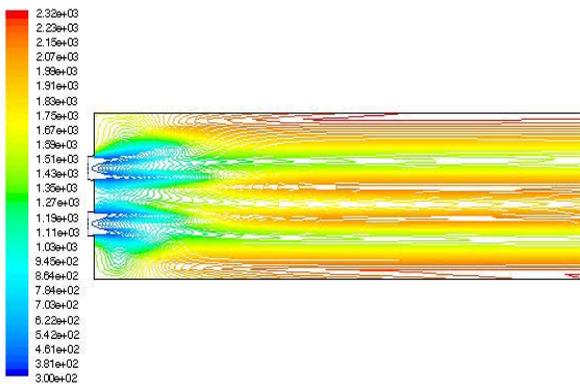


Fig. 13. The temperature fields in section plane with maximum heat density for case C4.

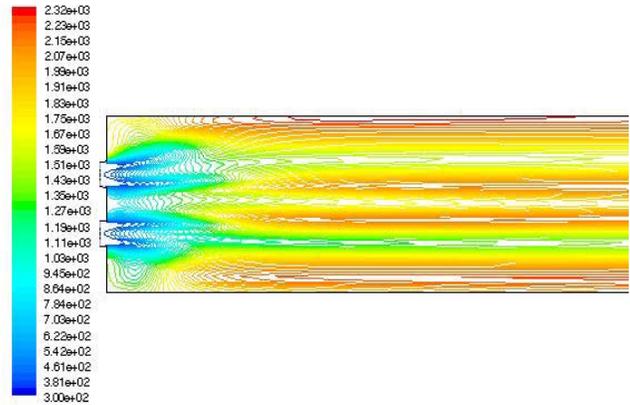


Fig. 14. The temperature fields in section plane with maximum heat density for case C8.

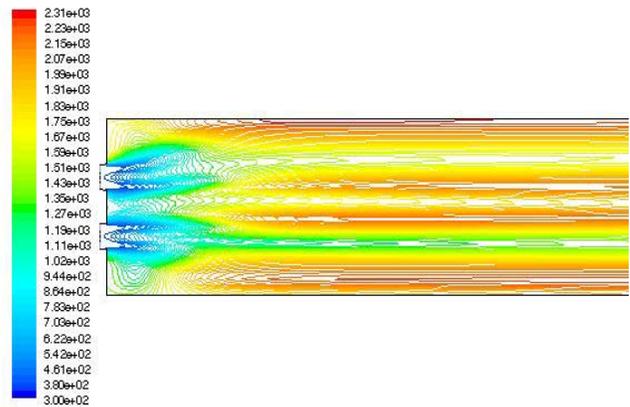


Fig. 15. The temperature fields in section plane with maximum heat density for case C10.

The numerical results of simulation for  $|\Delta T_{\max}|$  are given in Table V and for the deviations of sizes in [%] in Table VI.

Table V. The numerical results of simulation for  $|\Delta T_{\max}|$ .

	$ \Delta T_{\max} $
wall 1	0.25
wall 2	7.20
wall 3	28.0
wall 4	10.0
wall 5	0.08
output section	3.50

Table VI. The numerical results of simulation for the deviations of sizes in [%].

Physical sizes	Deviations of sizes in [%]		
p[Pa]	$\nearrow$ 100 %	$P_{NW}$ [ $Q_N$ , $T_N$ ]	$\nearrow$ 41.3 %
$\rho$ [kg/m <sup>3</sup> ]	$\nearrow$ 4.52 %		$\searrow$ 0.5%
T[K]	$\searrow$ 3.62 %		$\nearrow$ 0.55 %

v[m/s]	↗ 48.75 %		↗ 25.62 %
CH <sub>4</sub>	↗ 14 %		↘ 1.1%
CO <sub>2</sub>	↘ 3.14 %		↗ 0.47 %
O <sub>2</sub>	↗ 25 %		↘ 5%
N <sub>2</sub>	↘ 1 %		↗ 0.07 %
H <sub>2</sub> O	↘ 3.12 %		↗ 0.5 %

Fig. 17. The percentage of the variation for the maximum temperature of: T<sub>2</sub>, T<sub>3</sub>, T<sub>4</sub>, T<sub>e</sub> in function of Q<sub>1ef</sub>.

Fig. 18. The percentage of the variation for the maximum temperature of: T<sub>1</sub>, T<sub>5</sub> in function of Q<sub>1ef</sub>

To make the study of the temperature variation as consequence of the flow methane volume injected in burner Br1, is take as reference the correspondent sizes of the nominal regime of work, (case C6).

The plots have the values of physical sizes in percentage. Abscissae is the flow volume of methane injected in burner Br1 and the ordinate is the effective temperature. Considering the values from Table II and Table IV were traced using Maple 11 programme the variation of medium values in Fig. 16 and maximum values in Fig. 17 and Fig. 18.

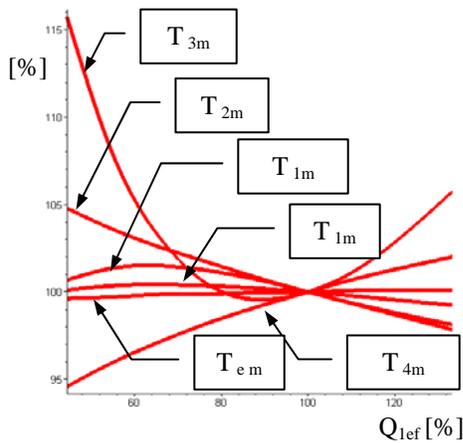
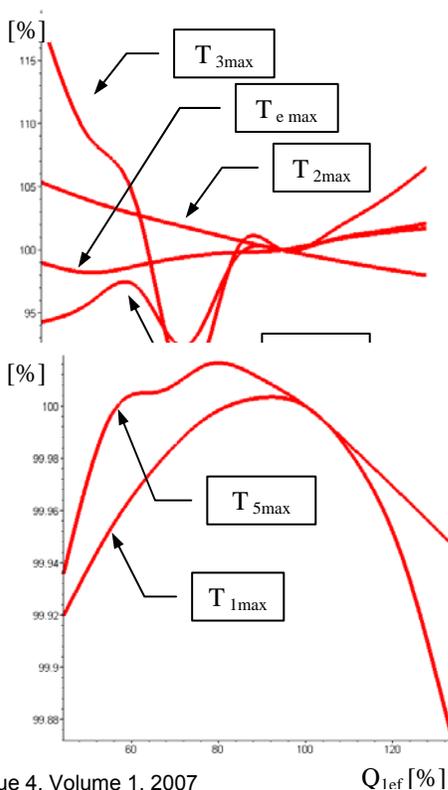


Fig. 16. The percentage of the variation for the medium temperature in function of Q<sub>1ef</sub>.



### III. CONCLUSIONS

The analysis with the Finite Element Method provides the turbulent flow regime with the internal and external recirculation zones, for fuel mixture injected into furnace and for the emitted gases resulted from the oxidation reaction of methane (Fig. 6).

With increasing the flow volume of methane injected, the maximum values of temperature field increases and moving to output section (Fig. 13 to Fig. 16).

Analyzing from Table IV the values of the modulus deviation for the maximum temperature, given in table VI, we can make the following remark:

- the bigger deviation of temperature in order, is on the frontal walls (the wall with 3, respective 28 %), the next on the output sections 3, with 5 % and the last on lateral the walls, (the wall 1, with 0.25 %).

In Table VI are given the percentage deviations in relation with the point of the nominal regime calculated as integral volume average into 3D domain of flow starting to values from Table III.

The p, v, O<sub>2</sub> and CH<sub>4</sub> have height deviations considering as reference their values from nominal point of work P<sub>NW</sub> [Q<sub>N</sub>, T<sub>N</sub>] and in opposition side the physical sizes ρ, T, CO<sub>2</sub>, N<sub>2</sub> and H<sub>2</sub>O have small deviations under Δ < 4.5 %.

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